

# Eddy-Mean Flow Interactions in Western Boundary Current Jets

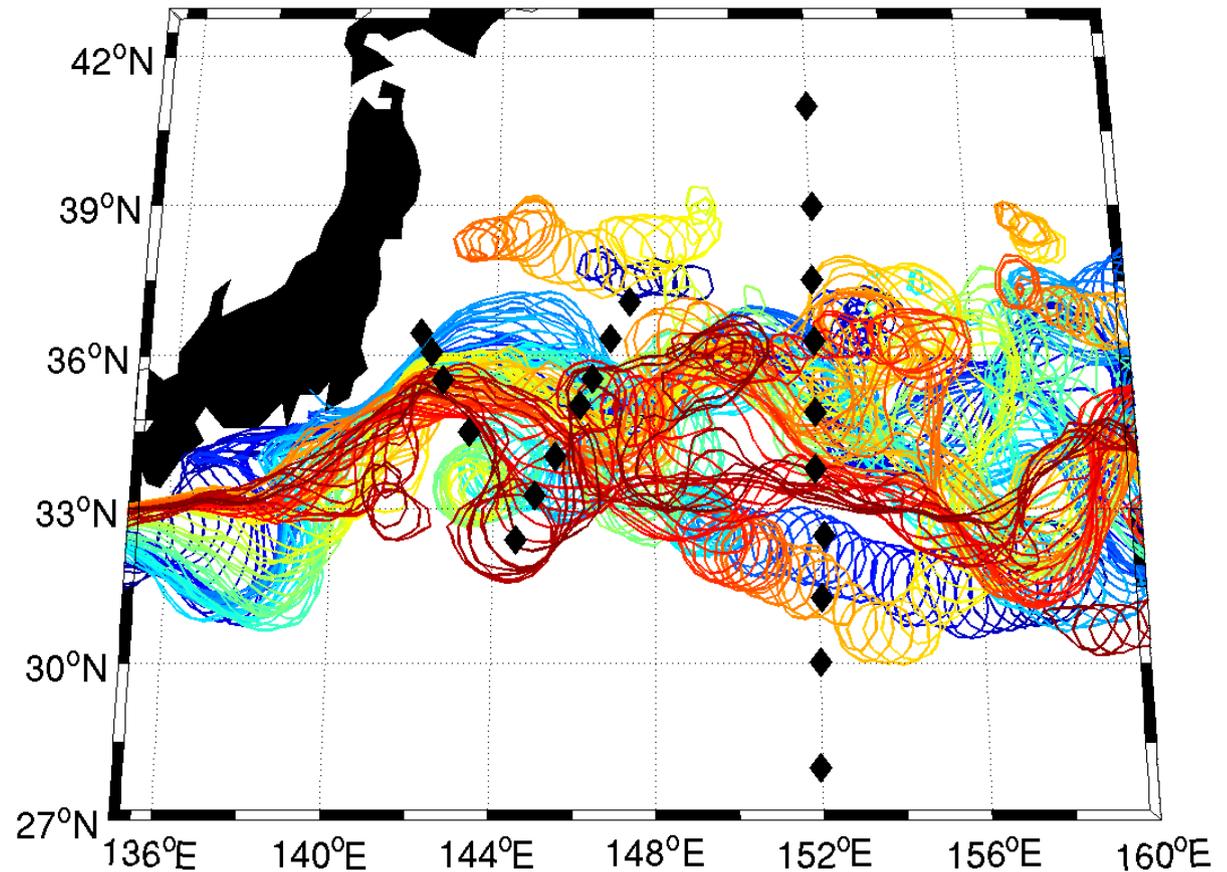
*An observationally driven theoretical study*

**Stephanie Waterman**  
MIT-WHOI Joint Program

**Steven Jayne**  
Woods Hole Oceanographic  
Institution

**Nelson Hogg**  
Cornell University

**March 3, 2008**



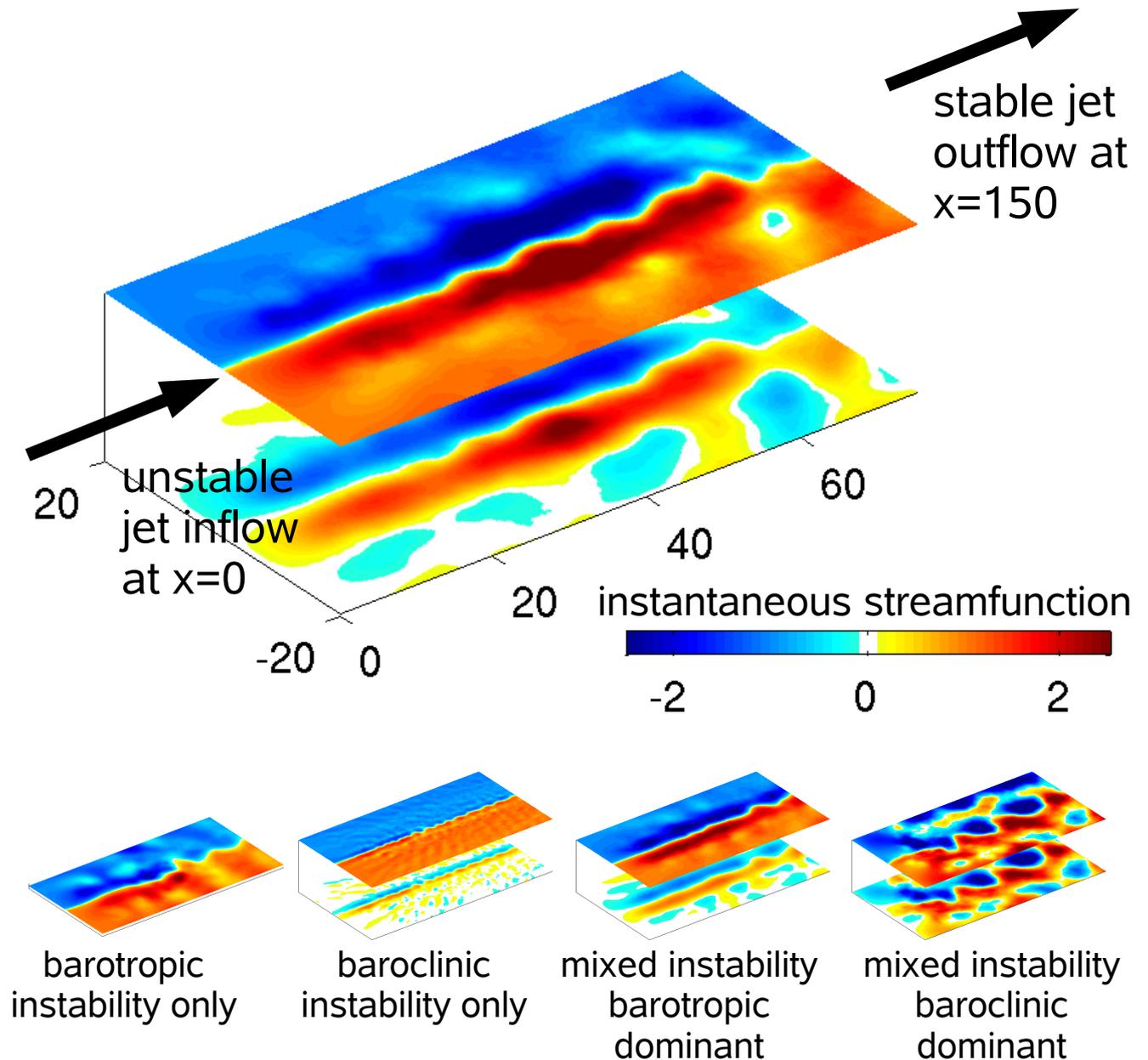
weekly snapshots of the 2.1 m SSH contour (proxy for the Kuroshio Extension jet axis) during the KESS observational period: May 2004 – June 2006



# The Model

a baroclinic, unstable, boundary-forced jet in an open domain

- QG
- fully non-linear
- 1 or 2-layer
- unstable jet inflow
- insensitive to outflow condition
- sponge layers on all boundaries to model “open ocean”
- posed in terms of time-mean and deviation from time-mean state



# The Model

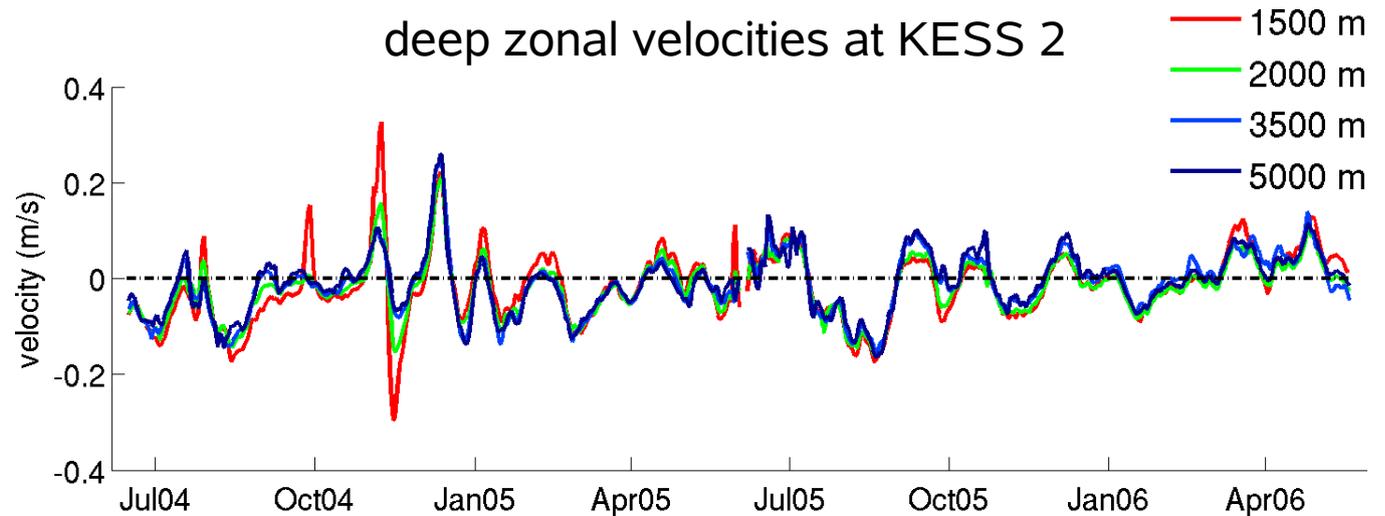
simplifications we employ and the physics we retain are appropriate to the Kuroshio Extension system

- weakly depth-dependent below the thermocline
- subject to mixed instability
- strongly nonlinear

# The Model

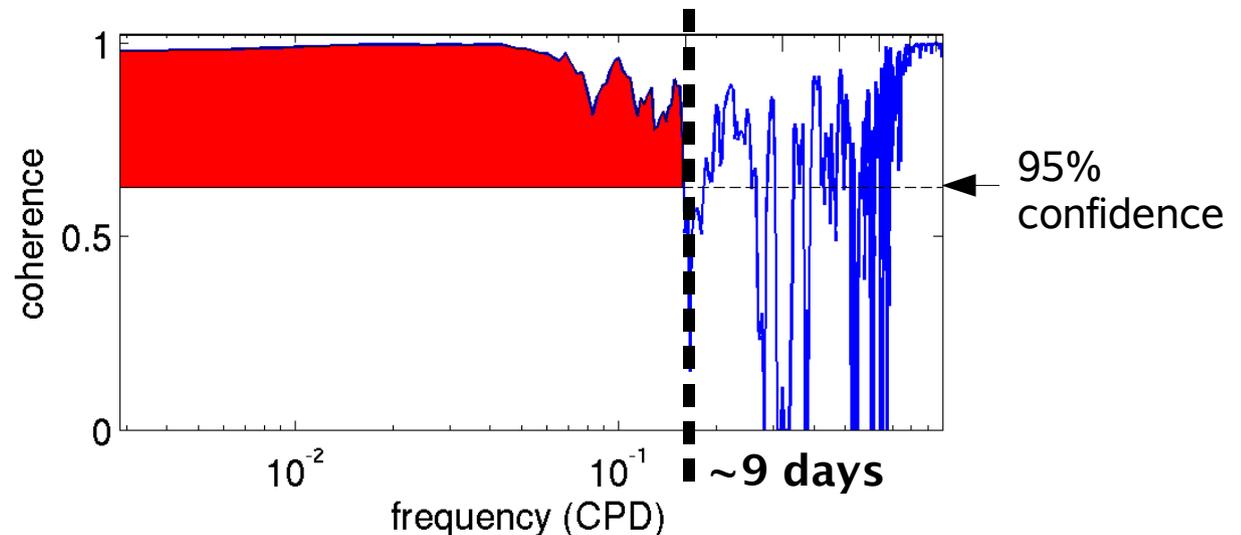
simplifications we employ and the physics we retain are appropriate to the Kuroshio Extension system

- weakly depth-dependent below the thermocline



- subject to mixed instability

vertical coherence between records at 1500 m and 5000 m



- strongly nonlinear

# The Model

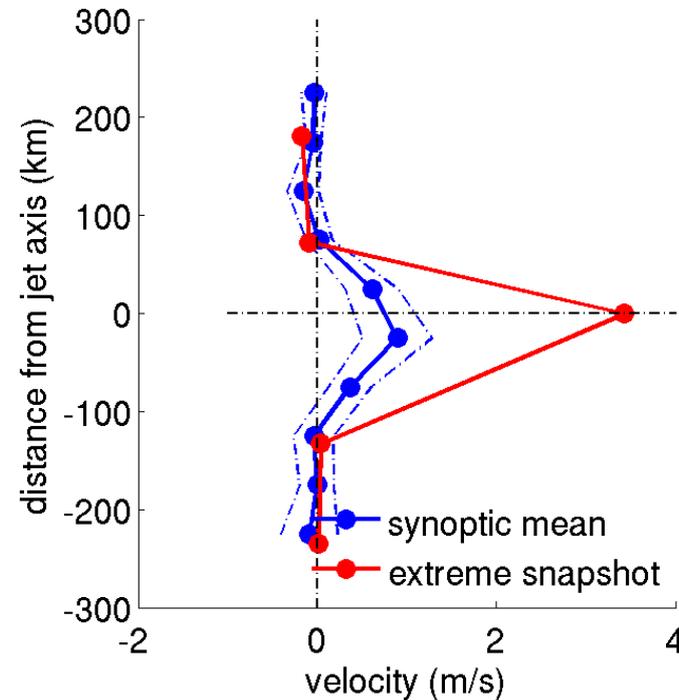
simplifications we employ and the physics we retain are appropriate to the Kuroshio Extension system

- weakly depth-dependent below the thermocline

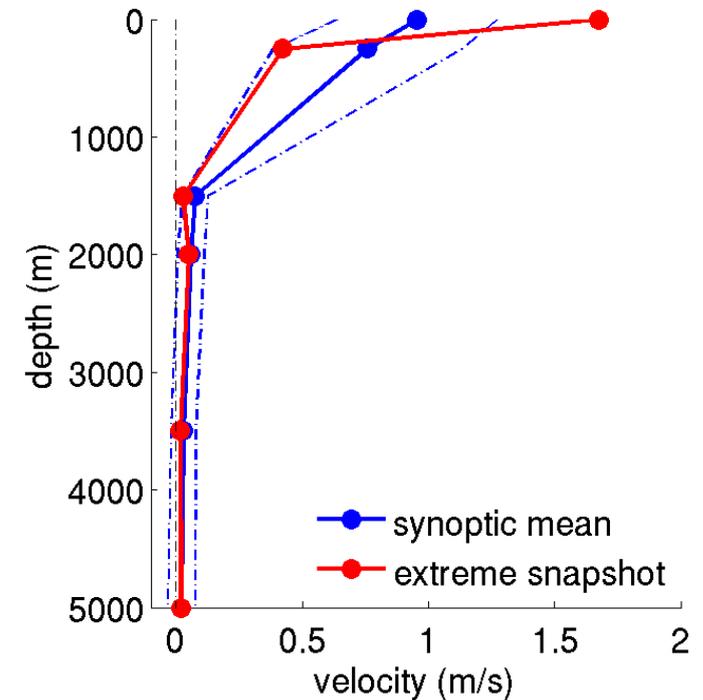
- **subject to mixed instability**

- strongly nonlinear

horizontal shear at 250 m



vertical shear



$$\left. \frac{\partial^2 U}{\partial y^2} \right|_{\text{mean}} \sim 2 \times 10^{-5} \frac{1}{\text{ms}}$$

$$\left. \frac{\partial^2 U}{\partial y^2} \right|_{\text{extreme}} \sim 7 \times 10^{-5} \frac{1}{\text{ms}}$$

$$\left( \text{vs. } \beta \sim 2 \times 10^{-11} \frac{1}{\text{ms}} \right)$$

$$\Delta U|_{\text{mean}} \sim 0.8 \frac{\text{m}}{\text{s}}$$

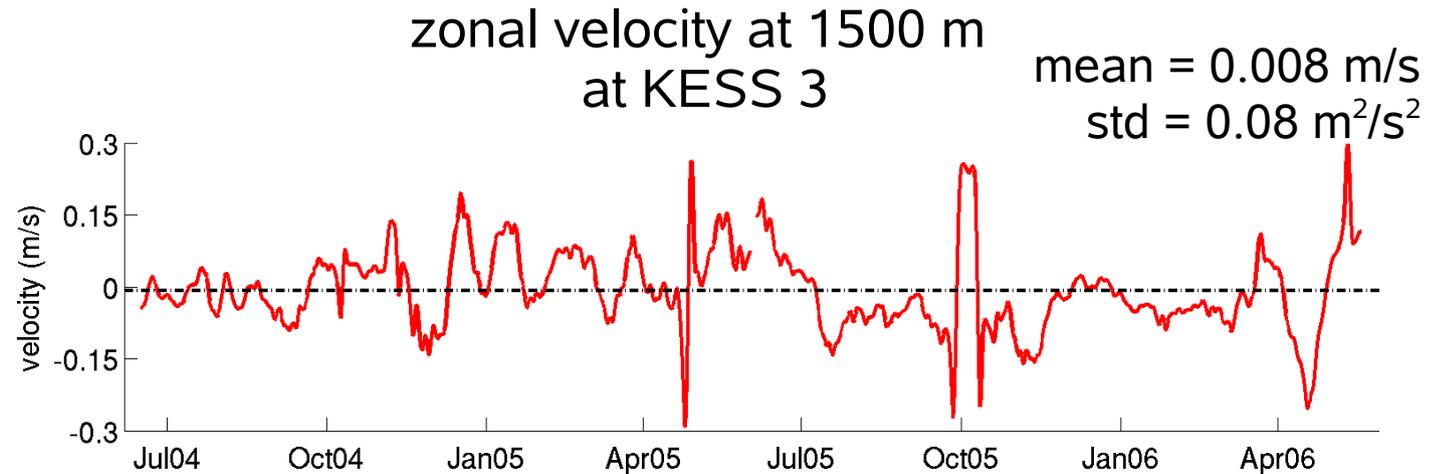
$$\Delta U|_{\text{extreme}} \sim 1.0 \frac{\text{m}}{\text{s}}$$

$$\left( \text{vs. } \Delta U|_{\text{critical}} \sim 0.1 \frac{\text{m}}{\text{s}} \right)$$

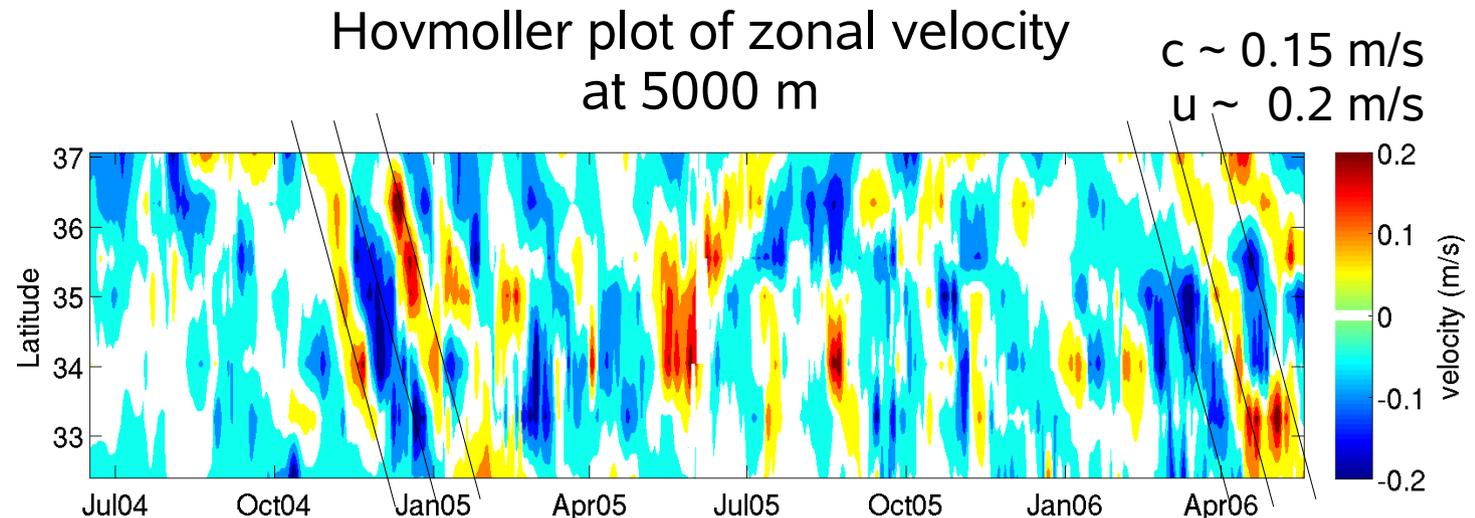
# The Model

simplifications we employ and the physics we retain are appropriate to the Kuroshio Extension system

- weakly depth-dependent below the thermocline



- subject to mixed instability



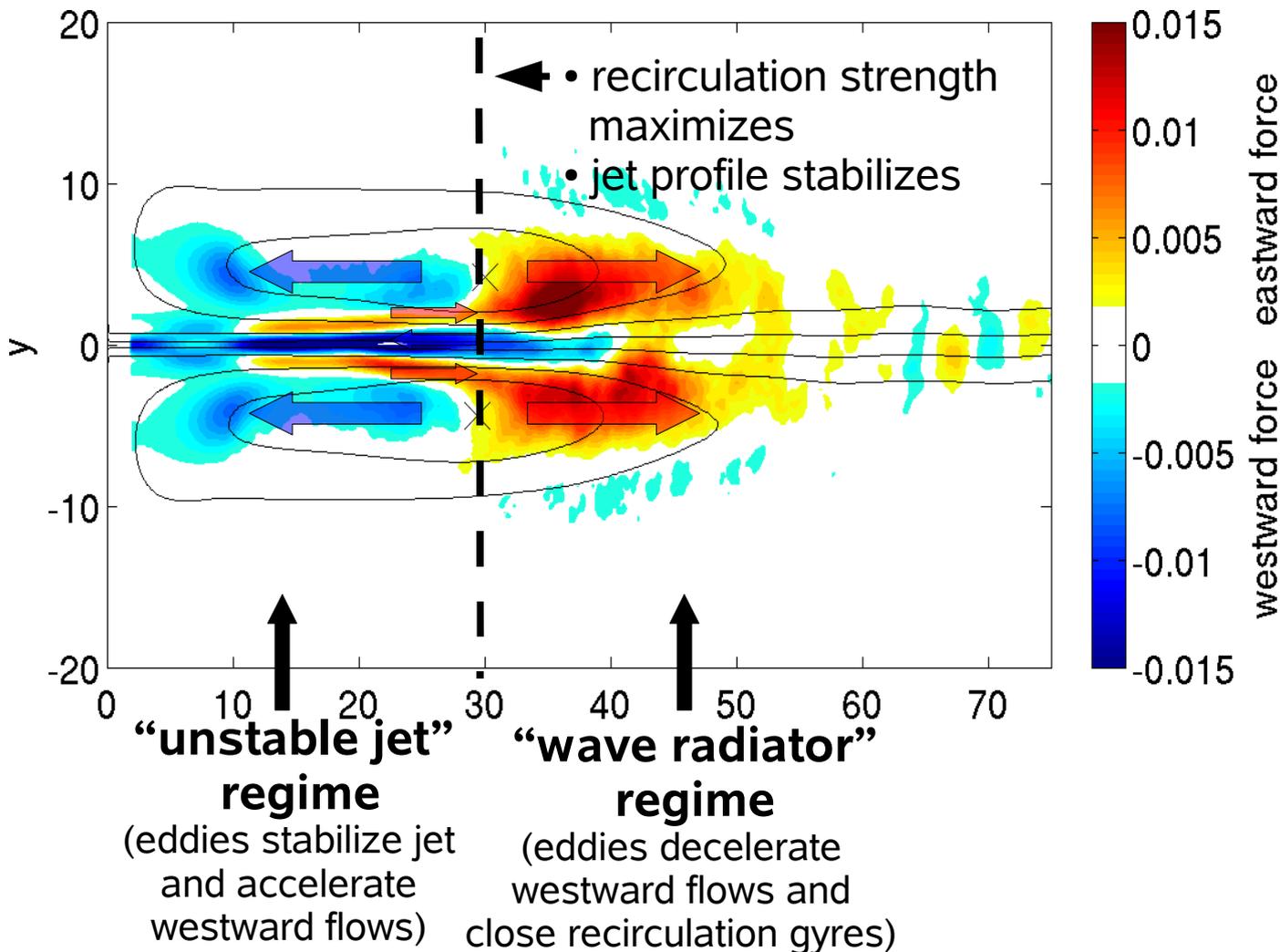
- **strongly nonlinear**

# Model Results:

eddies play a critical role in the downstream evolution of the jet through:

1. **stabilizing** the jet
2. **driving** the time-mean recirculations

## The Effective “Eddy Force”

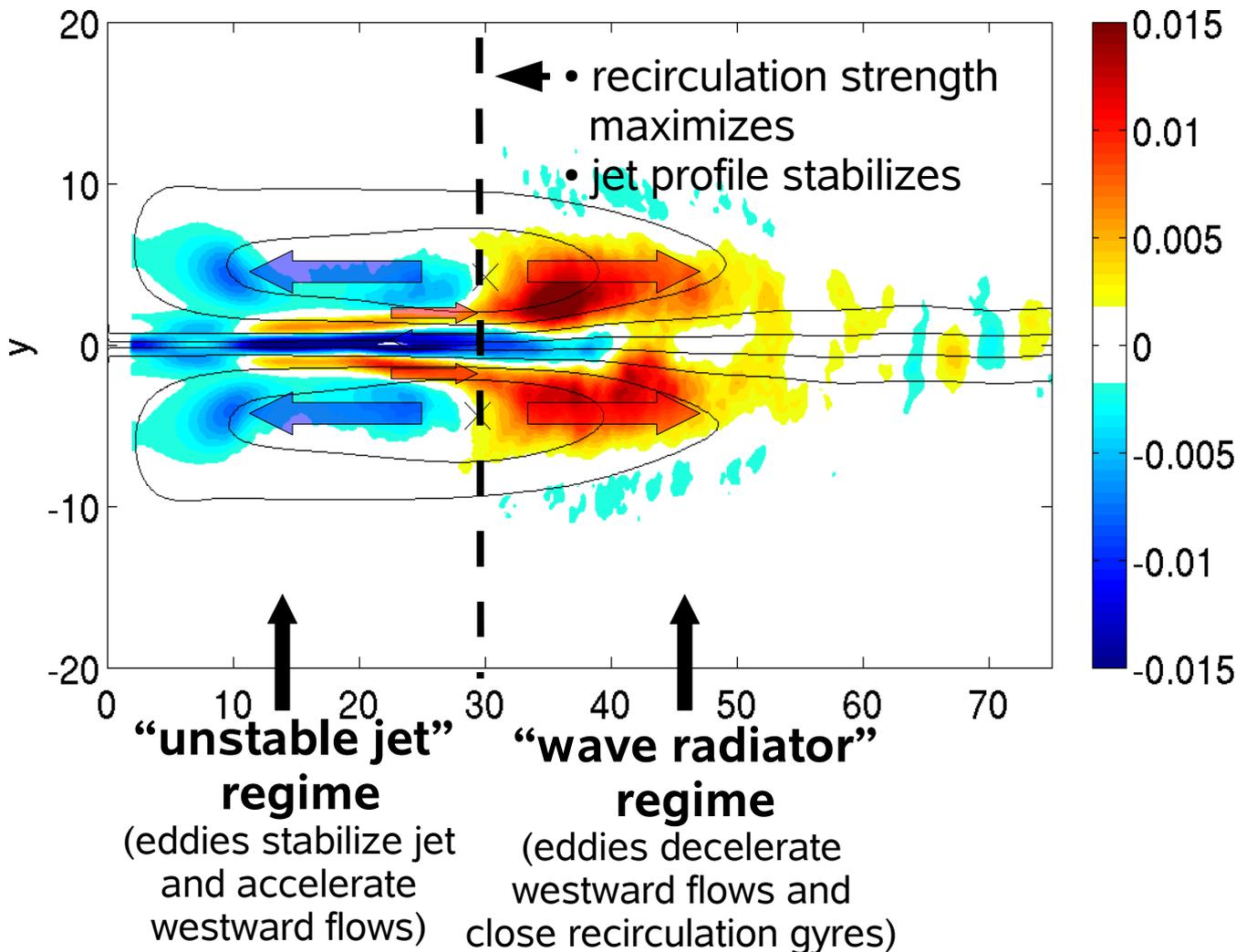


# Model Results:

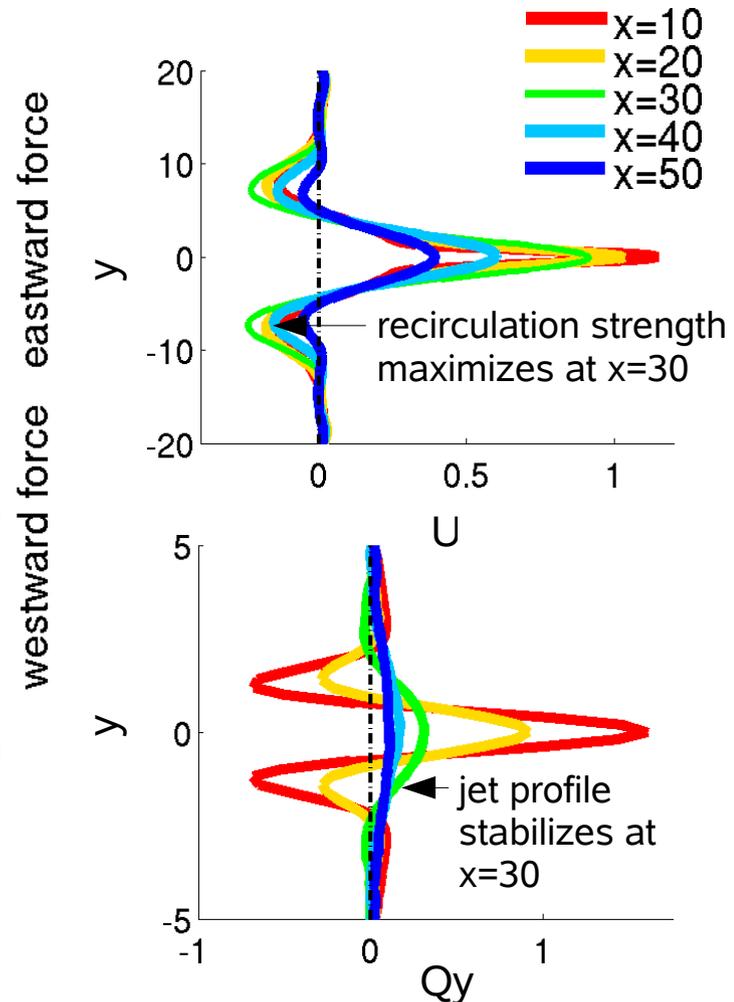
eddies play a critical role in the downstream evolution of the jet through:

1. **stabilizing** the jet
2. **driving** the time-mean recirculations

## The Effective “Eddy Force”



downstream evolution of U (top) and  $Q_y$  (bottom)

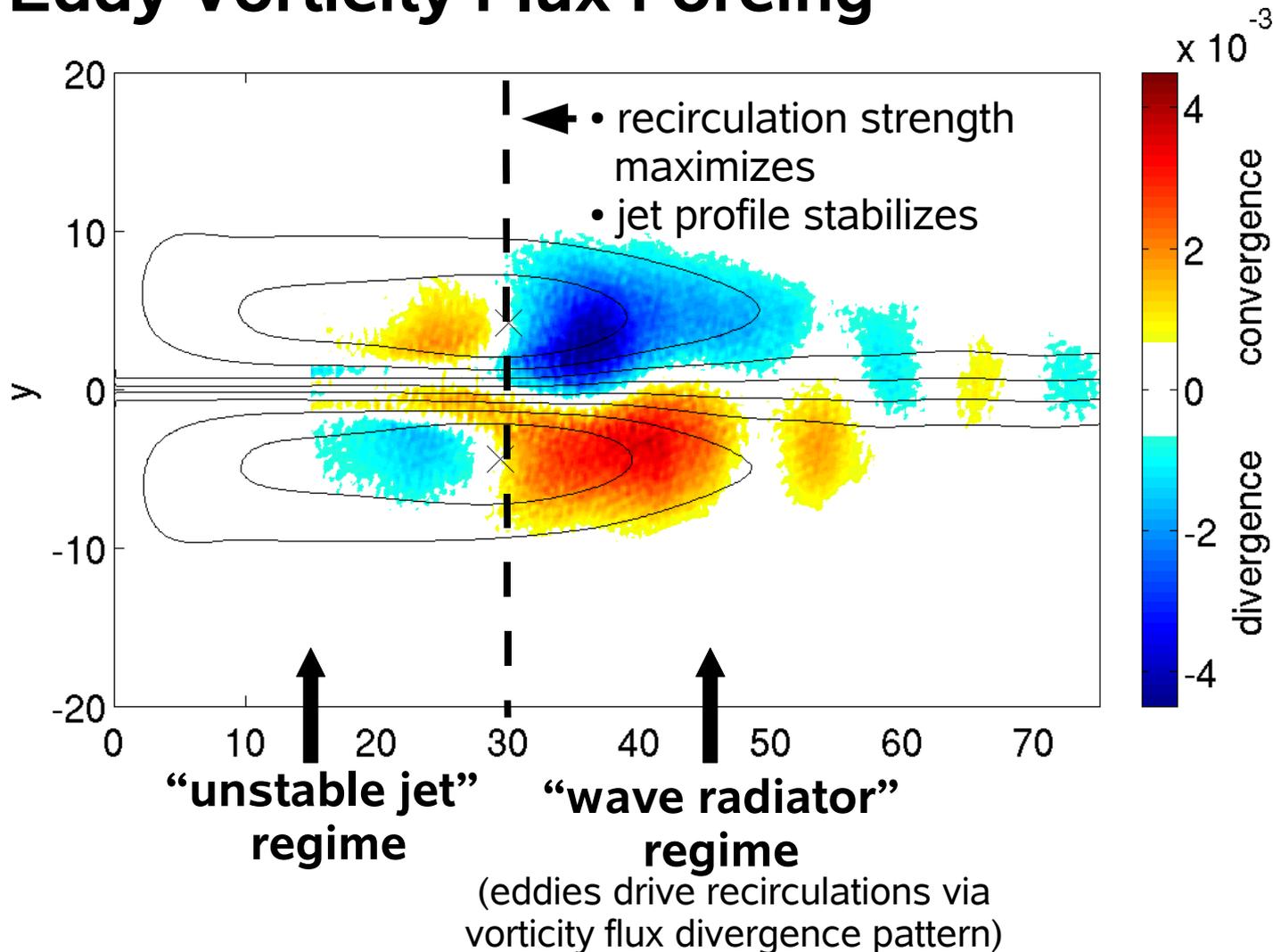


# Model Results:

eddies play a critical role in the downstream evolution of the jet through:

1. **stabilizing** the jet
2. **driving** the time-mean recirculations

## Eddy Vorticity Flux Forcing

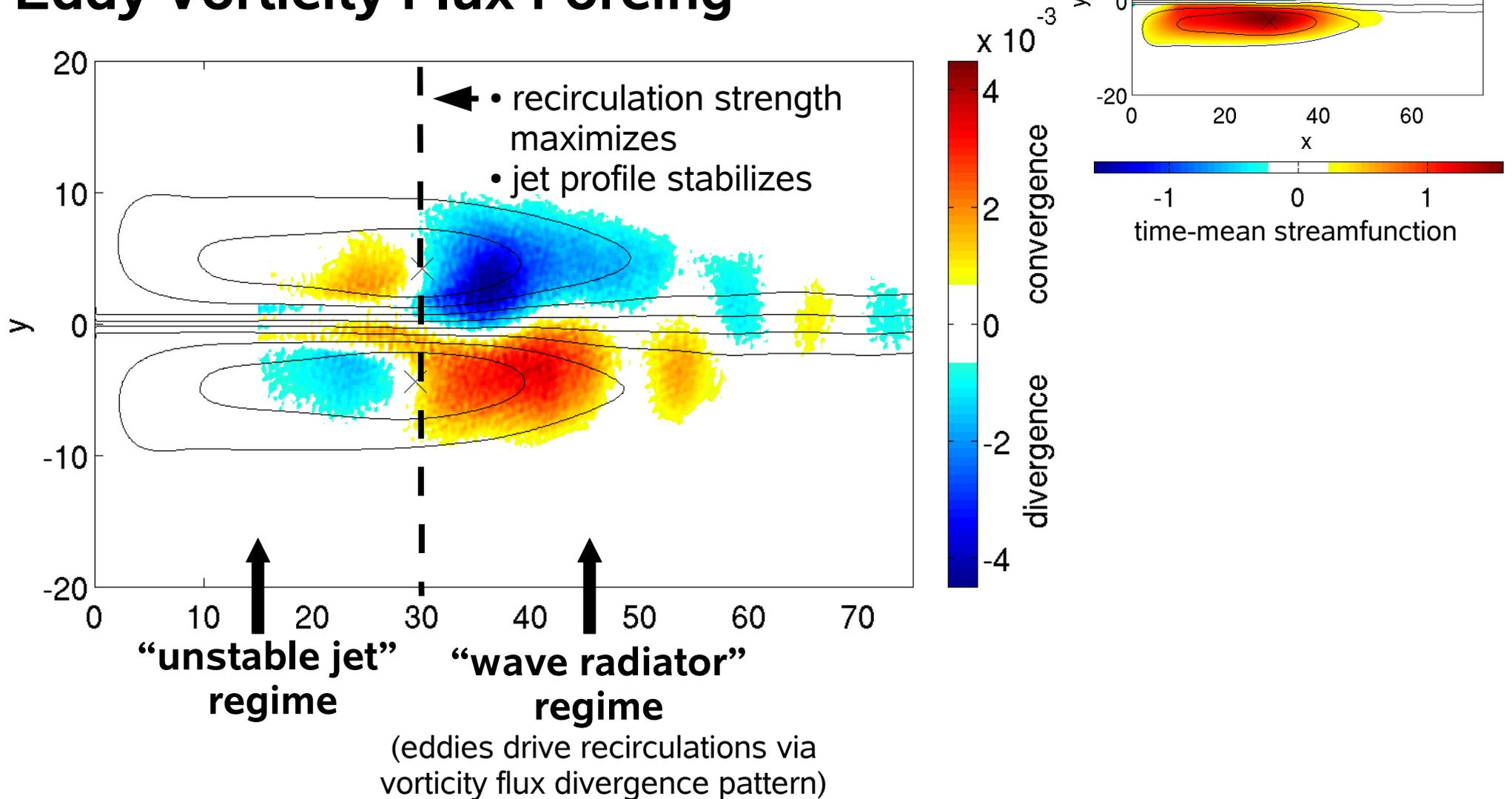


# Model Results:

eddies play a critical role in the downstream evolution of the jet through:

1. **stabilizing** the jet
2. **driving** the time-mean recirculations

## Eddy Vorticity Flux Forcing

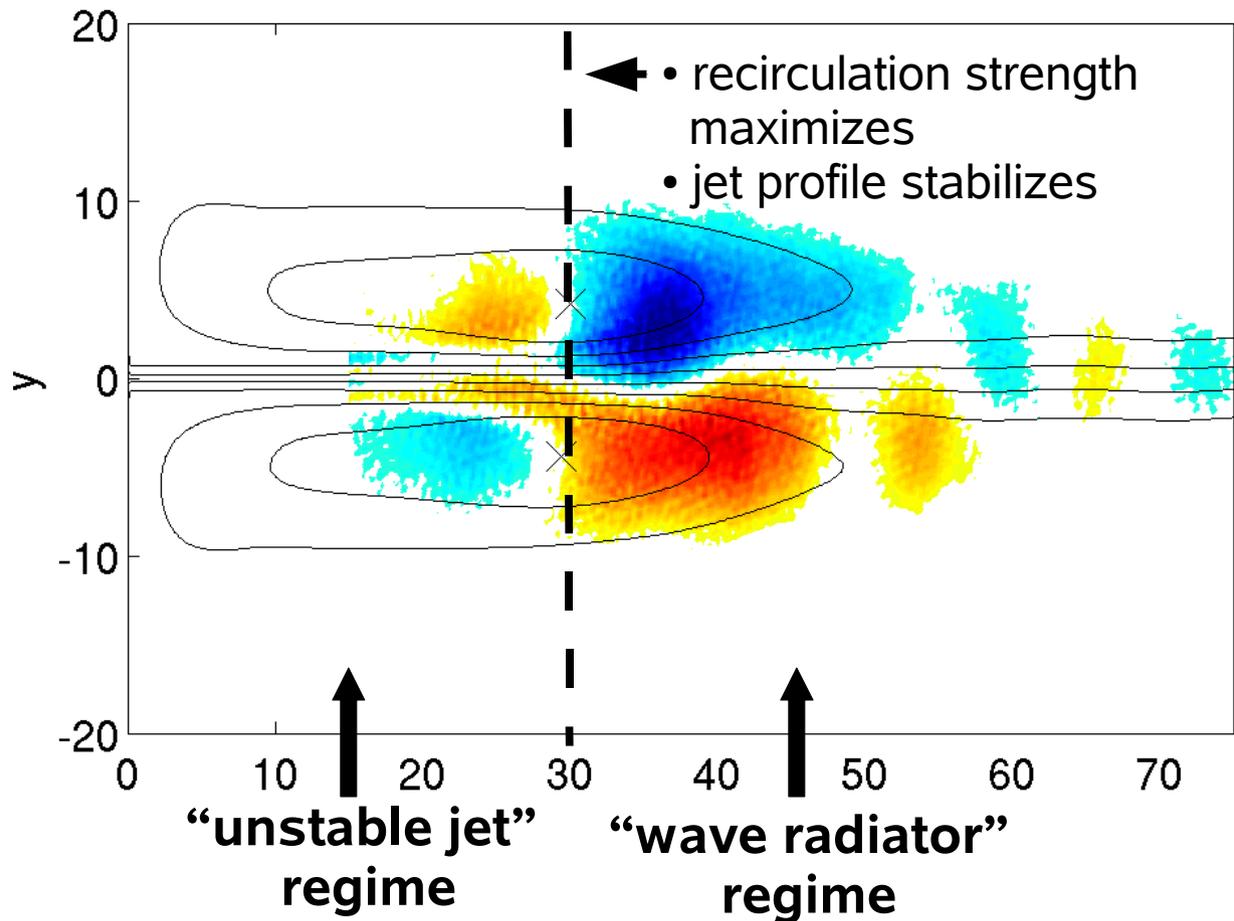


# Model Results:

eddies play a critical role in the downstream evolution of the jet through:

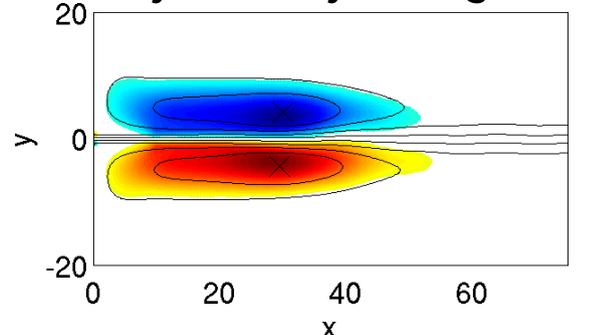
1. **stabilizing** the jet
2. **driving** the time-mean recirculations

## Eddy Vorticity Flux Forcing



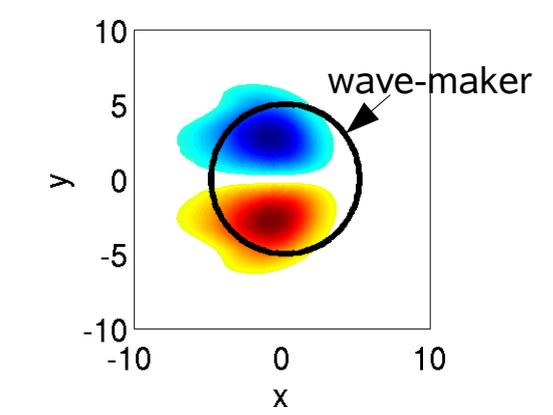
(eddies drive recirculations via vorticity flux divergence pattern)

time-mean circulation forced by the eddy forcing



time-mean streamfunction

eddy vorticity flux forcing from a localized wave source



vorticity flux divergence

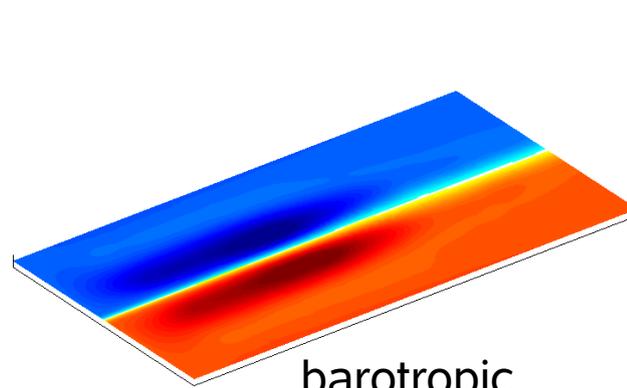
# Model Results:

the mechanism is unchanged by the addition of baroclinicity and/or baroclinic instability

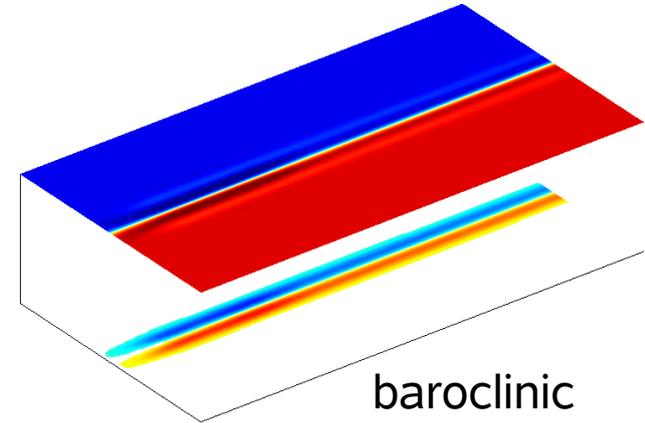
- regardless of unstable configuration - eddies drive recirculations

- baroclinic instability postpones (in  $x$ ) the barotropic mechanism: it creates (or adds to) the barotropically unstable jet

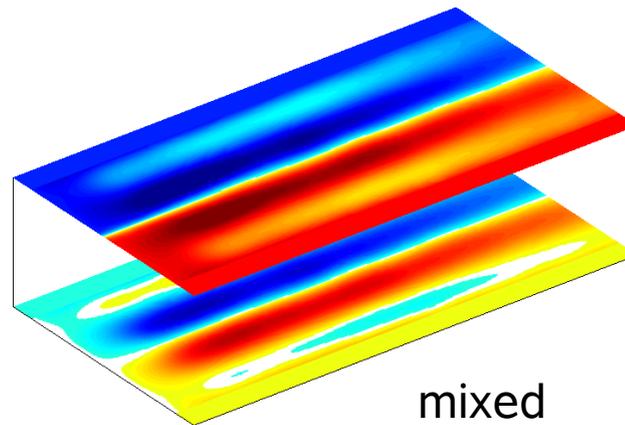
- new thickness fluxes reduce recirculation strength in the upper layer but drive lower layer recirculations



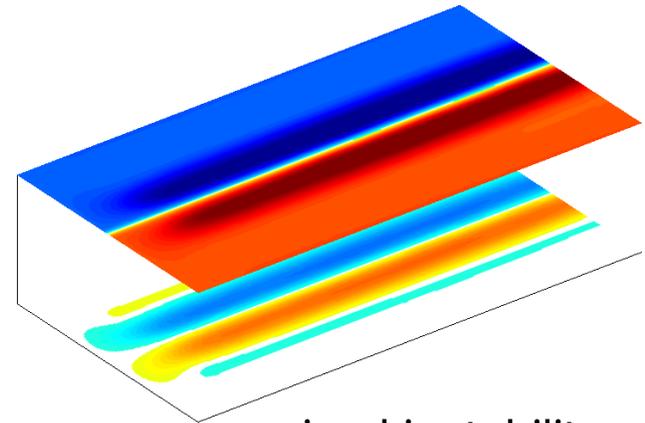
barotropic  
instability  
only



baroclinic  
instability  
only



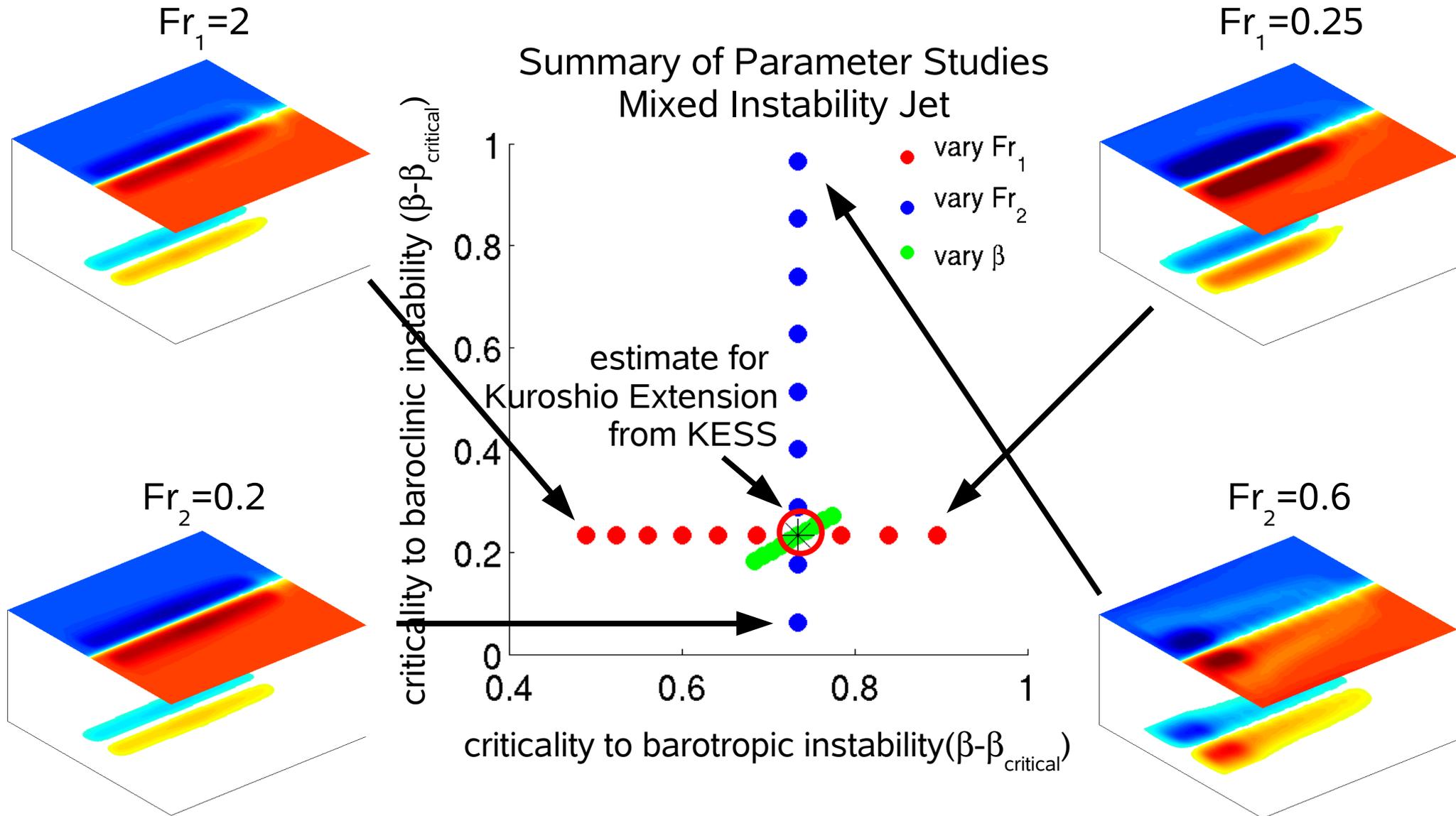
mixed  
instability  
(barotropic  
dominant)



mixed instability  
(baroclinic  
dominant)

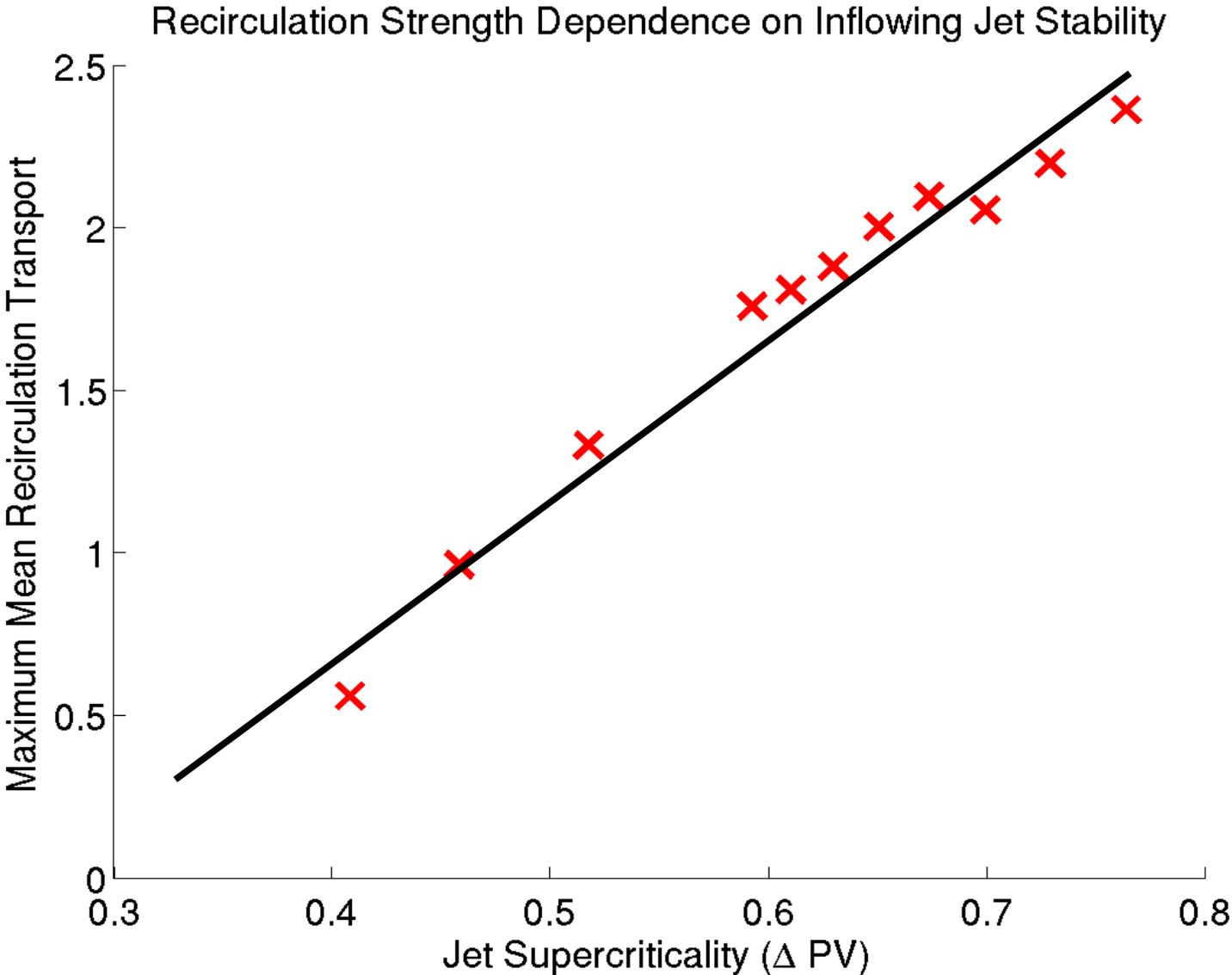
# Model Results:

eddy-driven time-mean circulation can be predicted empirically given the stability properties of the upstream jet



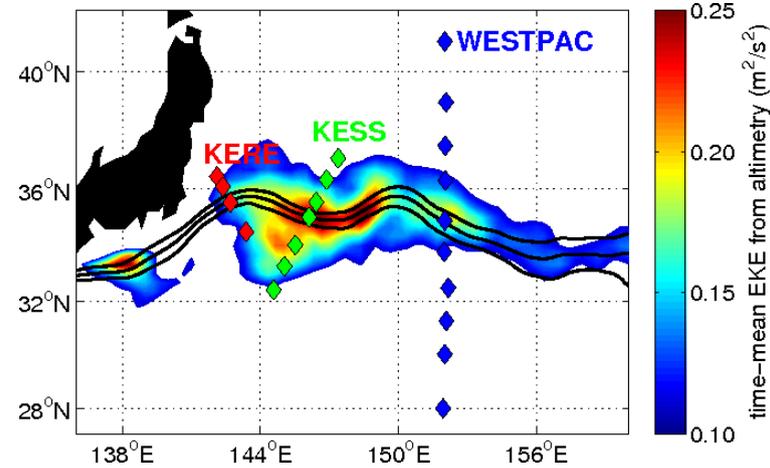
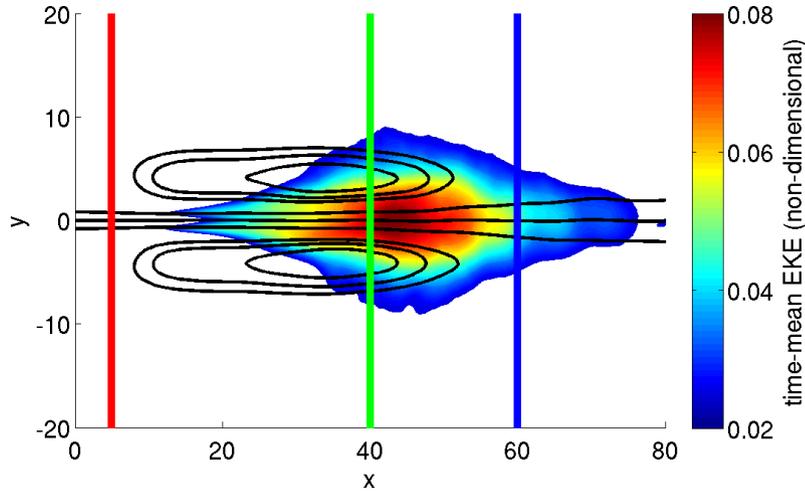
# Model Results:

eddy-driven time-mean circulation can be predicted empirically given the stability properties of the upstream jet



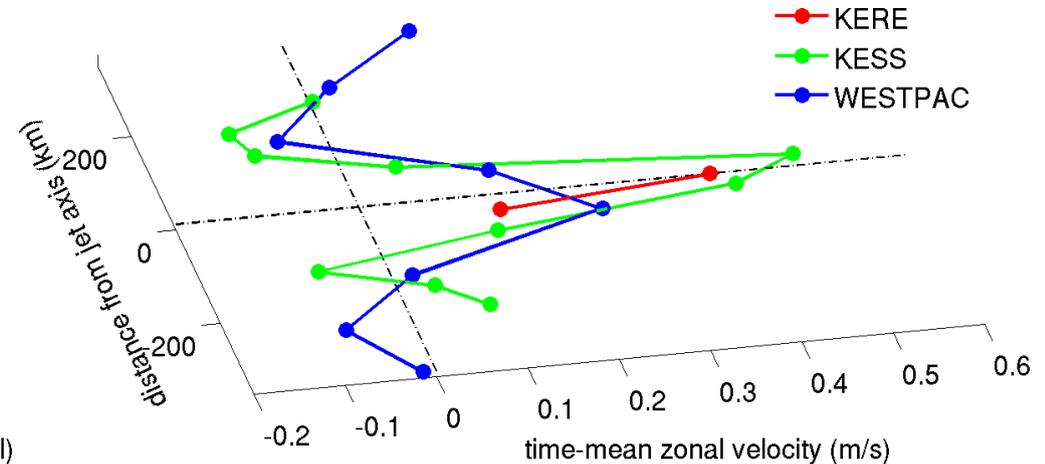
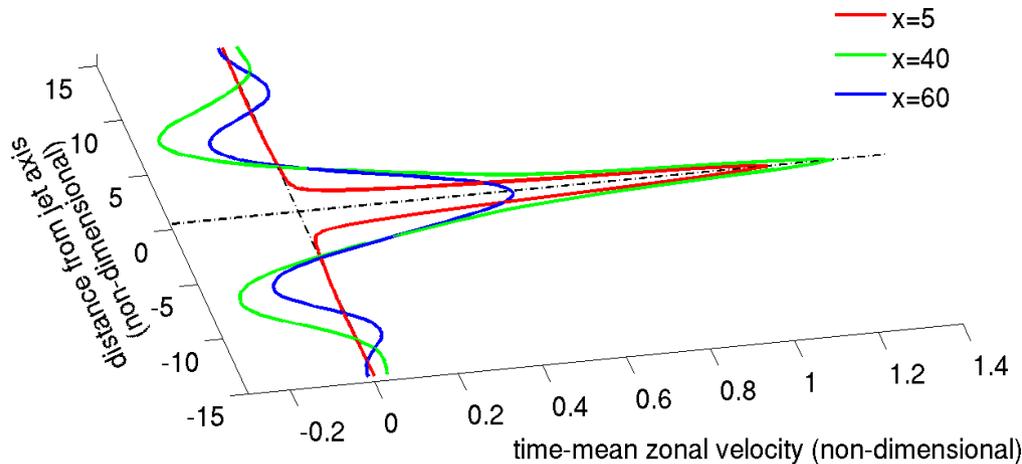
# Model – Observation Comparison:

consistencies in downstream development of **mean** and eddy properties suggest model has relevance to oceanic system



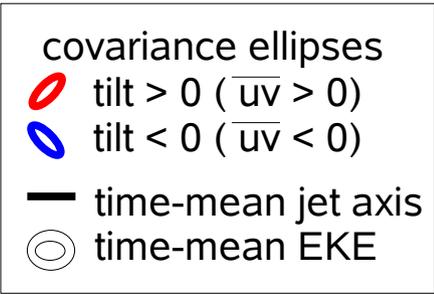
time-mean jet structure in upper layer  
model run with Kuroshio-like parameters

time-mean jet structure at 500 m  
observed

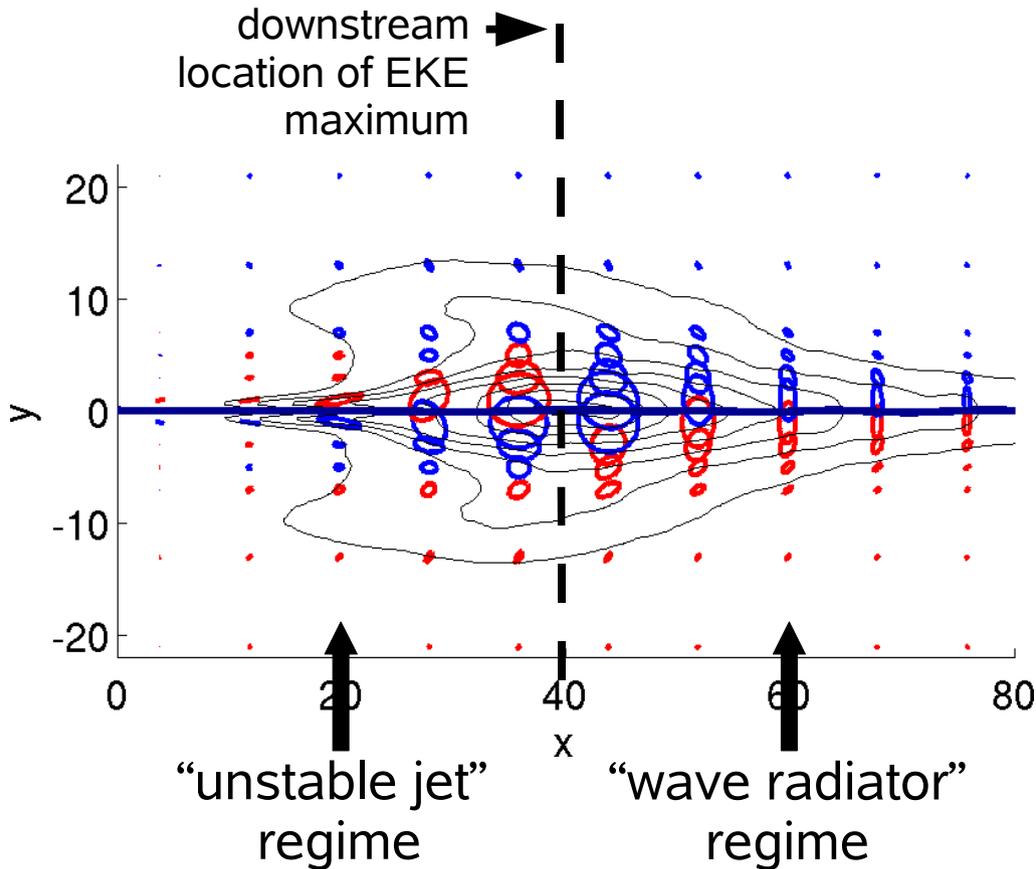


# Model – Observation Comparison:

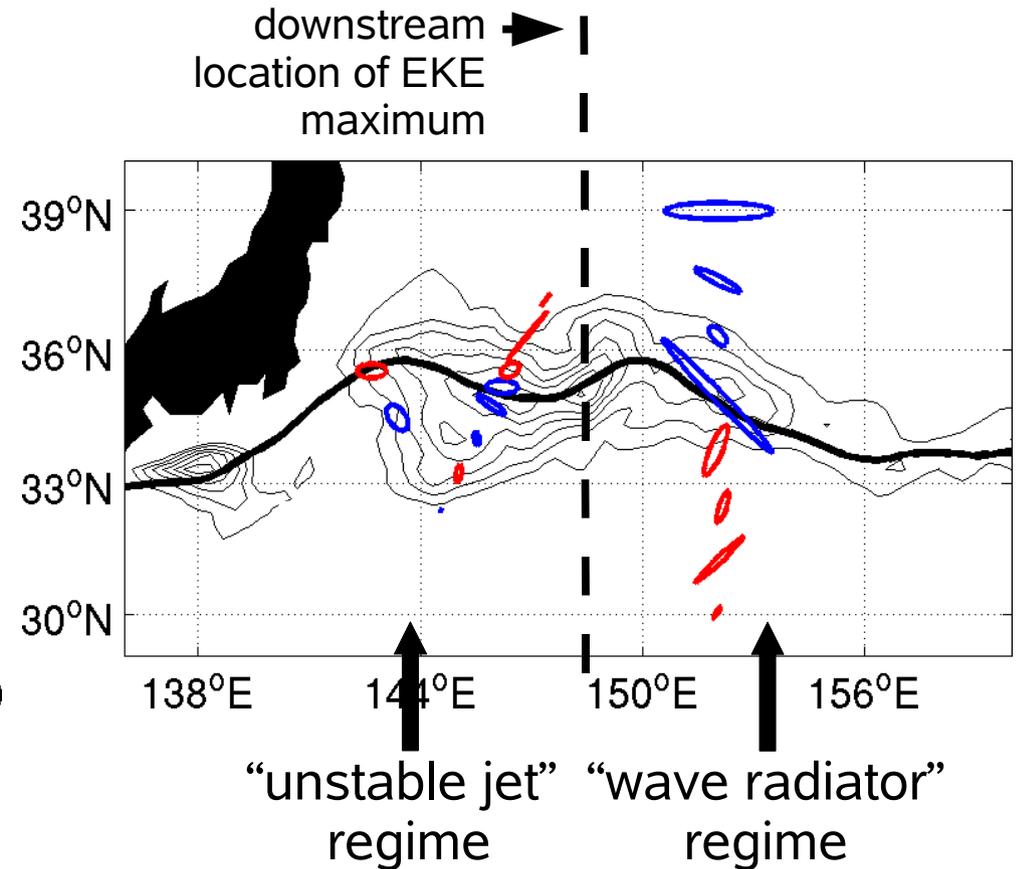
consistencies in downstream development of mean and **eddy** properties suggest model has relevance to oceanic system



barotropic model with Kuroshio-like parameters



depth-averaged observations

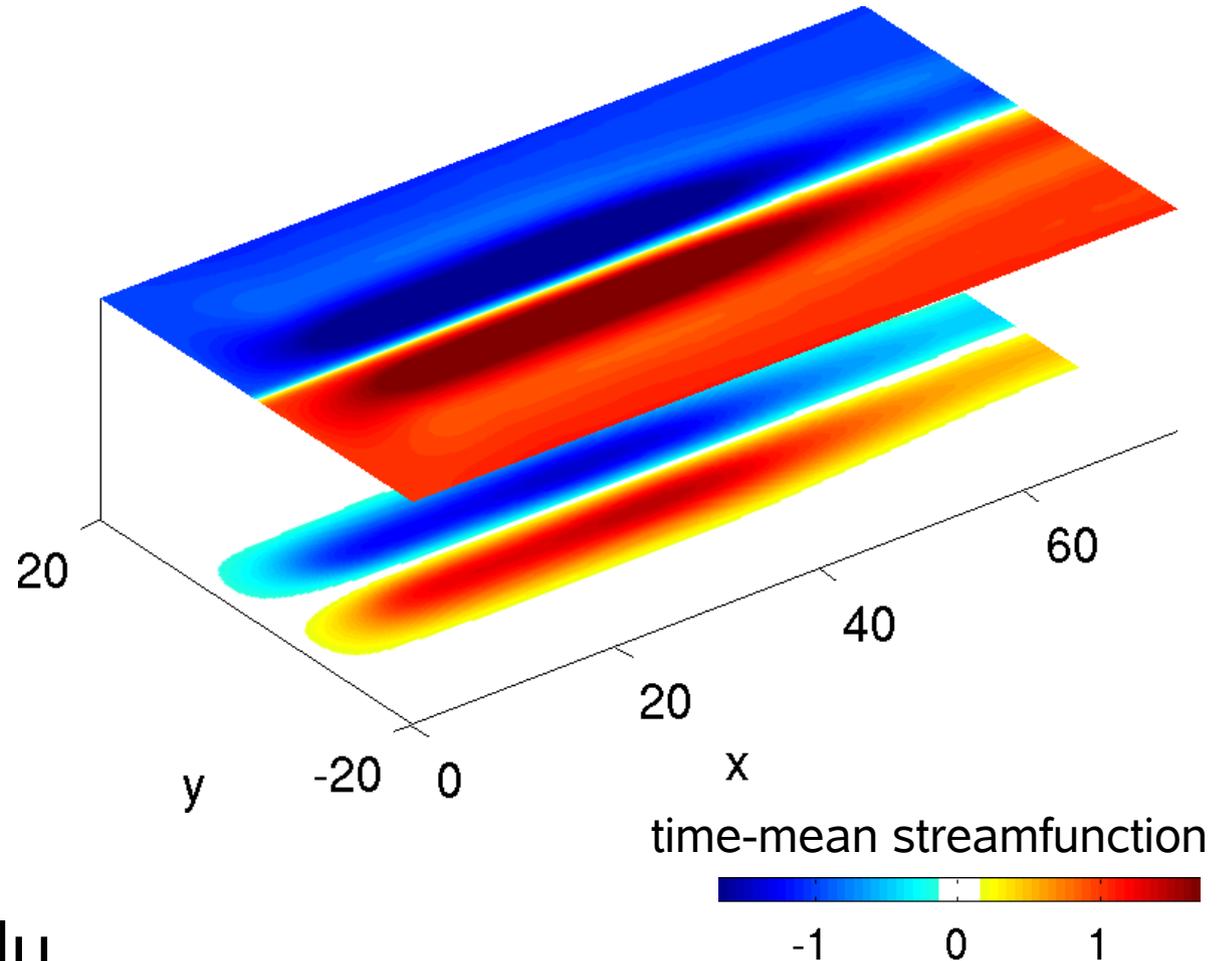


# In summary...

model teaches us the importance of eddy-mean flow interactions in the Kuroshio system

- eddies stabilize the jet
- eddies drive the time-mean recirculations
- zonal variation is important
- jet criticality determines mean recirculation properties
- model-observation consistencies suggest model has relevance to real oceanic system
- model in Kuroshio-like regime suggests Kuroshio is dominated by barotropic instability and eddies can drive recirculations of strength and extent consistent with observations

Time-mean circulation for a model run with Kuroshio Extension-like parameters ( $\beta=0.03$ ,  $Fr_1=1$ ,  $Fr_2=0.25$ )



More info: [snw@mit.edu](mailto:snw@mit.edu)